Exact controllability of the Schnakenberg model for Pattern Formation

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$$\begin{cases} u_t + \nabla \cdot (\mathbf{a}u) = \gamma(\alpha - u + u^2 v) + D_u \nabla^2 u \text{ in } Q \\ v_t + \nabla \cdot (\mathbf{a}v) = \gamma(\beta - u^2 v) + D_v \nabla^2 v \text{ in } Q \\ u(x, 0) = u_0(x), v(x, 0) = v_0(x) \text{ in } \Omega \\ \frac{\partial u}{\partial \nu}(x, t) = \frac{\partial v}{\partial \nu}(x, t) = 0 \text{ on } \partial\Omega \times (0, T) \end{cases}$$
(1)

- u, v: morphogen concentration.
- α, β : source density.

The Controllability Problem

Given a state (u^*, v^*) and fixed time T, Find α, β such that $u(T) = u^*, v(T) = v^*$, where u, v are solutions of (1)? **Theorem 1** Let $(u^*, v^*) \in H^2(\Omega) \times H^2(\Omega)$ be any steady-state solution to

$$\nabla \cdot (\mathbf{a}u^*) = \gamma(a^* - u^* + (u^*)^2 v^*) + D_u \nabla^2 u^* \text{ in } \Omega$$

$$\nabla \cdot (\mathbf{a}v^*) = \gamma(b^* - (u^*)^2 v^*) + D_v \nabla^2 v^* \text{ in } \Omega$$

$$\frac{\partial u^*}{\partial \nu} = \frac{\partial v^*}{\partial \nu} = 0 \text{ on } \partial \Omega$$
(2)

Then $\exists \delta > 0$ such that $\forall (u_0, v_0) \in H^1(\Omega) \times H^1(\Omega)$ satisfying

$$\|u^* - u_0\|_{\alpha_0(n)} + \|v^* - v_0\|_{H^1_0(\Omega)} \le \delta,$$

 $\exists \, \alpha, \beta \in \mathbb{R}^*_+$ such that

$$u^{\alpha,\beta}(T) = u^*, v^{\alpha,\beta}(T) = v^*.$$

Let $y = u - u^*, z = v - v^*$.

Our problem now is to prove the null controllability for

$$y_t + \nabla \cdot (\mathbf{a}y) = \gamma a + a_1 y + a_2 y^2 + a_3 z + a_4 y z + a_5 y^2 z + D_y \nabla^2 y$$

$$z_t + \nabla \cdot (\mathbf{a}z) = \gamma b + b_1 y + b_2 y^2 + b_3 z + b_4 y z + b_5 y^2 z + D_z \nabla^2 z$$

Initial condition: $y(0) = y_0 \equiv u_0 - u^*, z(0) = z_0 \equiv v_0 - v^*$

Linearized system:

$$y_{t} + \nabla \cdot (\mathbf{a}y) = \gamma a + y(\underbrace{a_{1} + a_{2}\xi}_{=\mu_{1}}) + z(\underbrace{a_{3} + a_{4}\xi + a_{5}\xi^{2}}_{=\mu_{2}}) + D_{y}\nabla^{2}y$$

$$z_{t} + \nabla \cdot (\mathbf{a}z) = \gamma b + y(\underbrace{b_{1} + b_{2}\xi}_{=\eta_{1}}) + z(\underbrace{b_{3} + b_{4}\xi + b_{5}\xi^{2}}_{=\eta_{2}}) + D_{z}\nabla^{2}z$$
(3)
where $\xi \in L^{\infty}(Q)$.

Lemma 1 $\exists a, b > 0$ such that $y_{\xi}(T) = z_{\xi}(T) = 0$, and

$$a^{2} + b^{2} \leq \frac{|\Omega|}{\gamma^{2}} \left(\|y_{0}\|^{2} + \|z_{0}\|^{2} \right) e^{\int_{0}^{T} \Phi(t) dt}$$
(4)

where

$$\Phi(t) = \max\{\|\mu_1\|_{L^{\infty}(\Omega)} + \frac{\|\eta_1\|_{\infty}}{2} + \frac{\|\mu_2\|_{\infty}}{2}, \|\eta_2\|_{\infty} + \frac{\|\eta_1\|_{\infty}}{2} + \frac{\|\mu_2\|_{\infty}}{2}\}.$$

Proof of Theorem 1. We use Kakutani's fixed point theorem with

$$K = \{\xi \in L^{\infty}(Q); \|\xi\|_{\infty} \le R\}$$

$$\Phi(\xi) = \{y_{\xi} \in L^{2}(Q) | (y_{\xi}, z_{\xi}) \text{ solution of (3)}, y_{\xi}(T) = z_{\xi}(T) = 0, a, b < M\}$$

Proof of Lemma 1.

Consider the optimal control problem

$$\min_{a,b} \left\{ \frac{a^2 + b^2}{2} + \frac{1}{\varepsilon} \int_{\Omega} \left(\|y(T)\|^2 + \|z(T)\|^2 \right) dt \right\}$$

subject to (3)

which has a unique solution $(a_{\varepsilon}, b_{\varepsilon})$ and $(y_{\varepsilon}, z_{\varepsilon})$.

By the Pontryagin maximum principle we have

$$a_{\varepsilon} = \gamma \int_{0}^{T} \int_{\Omega} p_{\varepsilon} \, dx dt, \quad b_{\varepsilon} = \gamma \int_{0}^{T} \int_{\Omega} q_{\varepsilon} \, dx dt,$$

where $(p_{\varepsilon}, q_{\varepsilon})$ is the solution to

$$-p_t - \mu_1 p - \eta_1 q + \mathbf{a} \cdot \nabla p - D_u \Delta p = 0,$$

$$-q_t - \mu_2 p - \eta_2 q + \mathbf{a} \cdot \nabla q - D_v \Delta q = 0,$$
 (5)

with final conditions

$$p(T) = -\frac{1}{\varepsilon} y_{\varepsilon}(T), \quad q(T) = -\frac{1}{\varepsilon} z_{\varepsilon}(T), \quad (6)$$

and boundary conditions

$$p\mathbf{a} \cdot \mathbf{n} - D_u \nabla p \cdot \mathbf{n} = 0, \quad q\mathbf{a} \cdot \mathbf{n} - D_v \nabla q \cdot \mathbf{n} = 0.$$

Then (3), (5) and (6) imply $\gamma^{2} \int_{0}^{T} \int_{\Omega} \left(p_{\varepsilon}^{2} + q_{\varepsilon}^{2} \right) dx dt + \frac{1}{\varepsilon} \int_{\Omega} \left(y_{\varepsilon}^{2}(T) + z_{\varepsilon}^{2}(T) \right) dx$ $= \int_{\Omega} \left(y_{0} p_{\varepsilon}(0) + z_{0} q_{\varepsilon}(0) \right) dx$ $\leq \frac{\gamma^{2}}{2C(T)} \int_{\Omega} \left(p_{\varepsilon}^{2}(0) + q_{\varepsilon}^{2}(0) \right) dx + \frac{C(T)}{2\gamma^{2}} \int_{\Omega} \left(y_{0}^{2} + z_{0}^{2} \right) dx$ $\leq \frac{\gamma^{2}}{2} \int_{0}^{T} \int_{\Omega} \left(p_{\varepsilon}^{2} + q_{\varepsilon}^{2} \right) dx dt + \frac{C(T)}{2\gamma^{2}} \int_{\Omega} \left(y_{0}^{2} + z_{0}^{2} \right) dx$

where

$$C(T) = \frac{1}{T} e^{\int_0^T \Phi(t) dt}.$$

Hence

$$\gamma^2 \int_0^T \int_\Omega \left(p_{\varepsilon}^2 + q_{\varepsilon}^2 \right) dx dt + \frac{2}{\varepsilon} \int_\Omega \left(y_{\varepsilon}^2(T) + z_{\varepsilon}^2(T) \right) dx \le \frac{C(T)}{\gamma^2} \int_\Omega \left(y_0^2 + z_0^2 \right) dx$$

On a subsequence we have

 $a_{\varepsilon} \to a, \quad b_{\varepsilon} \to b \quad \text{in } \mathbb{R},$ $y_{\varepsilon} \to y, \quad z_{\varepsilon} \to z \quad \text{weakly in } L^2(0,T;H^2(\Omega)) \cap W^{1,2}([0,T],L^2(\Omega))$ where (y,z) satisfy (3) and y(T) = 0, z(T) = 0.Moreover, a, b satisfy the estimate (4). This completes the proof of Lemma 1. The exact controllability for other Turing-type reaction-diffusion systems when initial condition is close to the desired state

- Gierer-Meinhardt Model
- Thomas Model
- Chemotaxis Model

Can this method be applied successfully for the reaction-diffusion system with quite general condition?

- 1. V. Barbu, Exact Controllability of the Superlinear Heat Equation, Appl Math Optim 42:73-89(2000).
- 2. V. Barbu, Local controllability of the phase field system, Nonlinear Analysis 50(2002) 363-372.
- 3. S. Kakutani, A generalisation of Brouwers fixed point theorem, Duke Math. J., 8 (1941), 457459.
- Ladyzhenskaya, Solonnikov, Uralceva, Linear and Quasi-Linear Equations of Parabolic type, AMS, Providence, Rhodes Island 1968.
- 5. A. Madzvamuse, Time-stepping schemes for moving grid finite elements applied to reaction diffusion systems on fixed and growing domains,
- 6. J. Murray, Math Biology. Volume I, Springer 2003.

Thank you for your attention!