Sparse Grid Collocation for Uncertainty Quantification

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 $\label{eq:http://people.sc.fsu.edu/~jburkardt/presentations/...} \\ \hfill \hf$



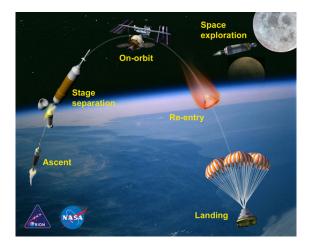
Sparse Grid Collocation for Uncertainty Quantification

Introduction

- Integrals Arising from Stochastic Problems
- Smolyak's Sparse Grid Definition
- Aspects of Sparse Grids
- Numerical Examples
- Conclusion



INTRO: Orion Space Vehicle Heat Shield





Kleb, Johnston, Uncertainty Analysis of Air Radiation for Lunar Return Shock Layers, AIAA 2008-6388.

INTRO: Wing Load Test

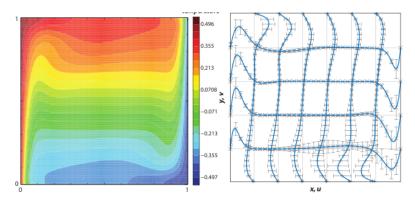




Babuska, Nobile, Tempone, Reliability of Computational Science, Numerical Methods for PDE's, Volume 23, Number 4, April 2007.

INTRO: Boussinesq Flow

Boussinesq flow, right hand side temperature is uncertain.



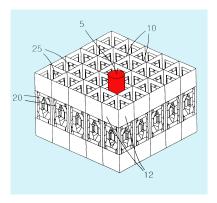
Temperature (left) and velocity (right) with uncertainty bars.

Le Maitre, Reagan, Najm, Ghanem, Knio, A stochastic projection method for fluid flow II, Journal of Computational Physics, Volume 181, pages 9-44, 2002.



INTRO: Fuel Rod Fretting

Nuclear fuel rods are held in a grid. The flow of coolant can set up vibrations. If resonance occurs, a rod suffers severe fretting.





Nam Dinh, CASL Validation Data: An Initial Review, INL/EXT-11-21017, Idaho National Laboratory, 2011. The preceding images suggest some of the situations in which a calculation must be carried out in the presence of uncertainty, error, or variation.

One can ignore the uncertainty, assume the data is exact, and compute an "exact answer".

One understands that this "exact answer" will not correspond to the physical situation actually encountered, in part because of the uncertainties we ignored; if we expect the actual behavior to be close to our calculation, that is simply a matter of hope.

However, a variety of techniques are available for making a reasonable model of uncertain input parameters, and then, for the output solution and related quantities of interest, computing the expected value, the variance, and other statistical quantities.



For a space vehicle subject to collisions with debris, I might begin by considering a quantity such as the particle size, which can have a continuous range of values. From observational records, I may choose a probability density function (PDF) that is a good model for the variability of this quantity.

We may also want to consider the velocity of the particle. This represents a second "dimension" of uncertainty. If I can assume the quantities are independent, then I can first construct a PDF for the velocities alone, and then combine the two PDF's to give probabilities about a particle of a given size and velocity.



Before we zoom off into higher dimensions, let's consider a simple case of uncertainty, and how we quantify it (turn it into a number).

I must pay \$5 to play a dice game, in which I roll two dice, and win as many dollars as the sum of the dice. Should I play?

The first answer is, I am uncertain of the outcome. I cannot predict what will happen on one game. There are really two independent stochastic parameters, ω_1 and ω_2 controlling the behavior of the separate dice. But let's simplify this to a single parameter ω that controls both.

Although I can't predict the result of any one game, I can consider all possible games, whose values to me range from \$2 to \$12. I can estimate my risk by imagining playing the game many many times, and averaging the winnings.



This turns out to the an estimate of the probability sum, which sums the product of each result with its probability to give me what is called the expected value.

$$E(\text{prize}) = \sum_{i=2}^{12} \text{prize}(i) p(i)$$

When our variable is continuous, we have instead

$$E(f) = \int_{\Omega} f(x(\omega)) \rho(\omega) \, d\omega$$

where $\rho(\omega)$ is the probability density function that says how likely each value of ω is.

When many separate uncertainties $\omega_1, \omega_2, \dots$ are involved, this becomes a multidimensional integral.



So, analyzing uncertainty requires estimating a high-dimensional integral.

Monte Carlo sampling produces a cheap estimate that is not very precise.

<u>Quadrature rules</u> spend effort making a model of the data, and then making an estimate of the integral whose accuracy can be controlled. To get similar accuracy in higher dimensions, we use product rules, but the cost of making the model breaks the budget (even 10 dimensions can be too many).

This talk considers a method (cheap and precise) for estimating integrals associated with some kinds of uncertain or stochastic problems.



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Stochastic problems are mathematical models of systems in which some data is uncertain, approximate, or nondeterministic.

Let us consider the Poisson diffusion equation:

 $-\nabla \cdot (a(\vec{x})\nabla u(\vec{x})) = f(\vec{x})$

 $\mathbf{a}(\vec{x})$ is the diffusivity, $\mathbf{f}(\vec{x})$ a source term.

A stochastic version might allow uncertainties in the diffusivity:

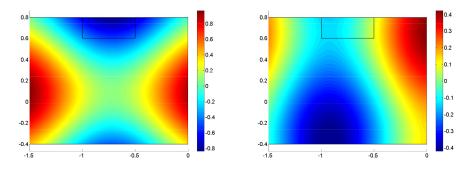
$$-\nabla \cdot (a(\vec{x}; \boldsymbol{\omega}) \nabla u(\vec{x}; \boldsymbol{\omega})) = f(\vec{x})$$

Here, ω represents the stochastic influence, and we must now write **u** with an implicit dependence on ω as well.



STOCHASTIC: Stochastic Diffusivity

Examples of a mildly varying stochastic diffusivity $a(x, y; \omega)$.



Log of realizations of the diffusivity $a(x, y; \omega)$ using uniform random variables ω (left) or Gaussian ω (right).



 $http://people.sc.fsu.edu/\sim jburkardt/m_src/stochastic_diffusion/stochastic_diffusion.html$

Since ω is unknown, our problem has a family of solutions, and any particular solution is actually not very important or meaningful, unless we know its associated probablity (PDF value, actually).

But what is important is to characterize the statistics of the solutions, such as the average behavior and variance from that average.

We could imagine addressing this question by solving every possible variation of the problem, and weighting it by its probability.

With each value of ω we associate a probability density function $\rho(\omega)$, and reformulate our problem to solve for the finite element coefficients of the expected value function $\mathbf{\bar{u}}(\vec{x})$.

The validity and usefulness of all our results depends on how carefully we choose a probabilistic model to combine with our deterministic system. $_{2}$



A finite element approach to the deterministic problem integrates the equation against various test functions $v_i(\vec{x})$:

$$\int_D a(\vec{x}) \nabla u(\vec{x}) \cdot \nabla v_i(\vec{x}) \, d\vec{x} = \int_D f(\vec{x}) v_i(\vec{x}) \, d\vec{x}$$

If we include an uncertain parameter ω , then we are trying to compute:

$$\int_D a(\vec{x}; \boldsymbol{\omega}) \, \nabla u(\vec{x}; \boldsymbol{\omega}) \cdot \nabla v_i(\vec{x}) \, d\vec{x} = \int_D f(\vec{x}) v_i(\vec{x}; \boldsymbol{\omega}) \, d\vec{x}$$

We don't actually know ω , so we don't really want to solve this problem; instead, we are interested in things like the expected value of the solution. To get this, we need to average over all possible values of ω :

$$\int_{\Omega} \int_{D} a(\vec{x};\omega) \nabla u(\vec{x};\omega) \cdot \nabla v_{i}(\vec{x}) d\vec{x} \rho(\omega) d\omega = \int_{\Omega} \int_{D} f(\vec{x}) v_{i}(\vec{x};\omega) d\vec{x} \rho(\omega) d\omega$$

Once we interchange the orders of integration, this looks very much like an ordinary finite element computation, an algebraic system A * c = f for the finite element coefficients of $\mathbf{\bar{u}}(\vec{x})$.

Evaluating A and f requires approximate integration over D and $\Omega.$



Solving the stochastic problem begins by approximating the probability space by some discretized *M*-dimensional subspace Ω^M . We are then faced with the problem of estimating these integrals.

Clearly, we can do a better job of estimating the influence of the stochastic parameters if we can make M large; on the other hand, the task of estimating integrals in M-dimensional space can quickly become intractable as the dimension increases.

We'll now concentrate on the problem of integral approximation, and thus replace our complicated stochastic diffusion integral by the model task of estimating integrals in an *M*-dimensional product region:

 $\int_{\Omega^M} f(x) \, dx$



One choice is the **Monte Carlo** method, which randomly samples the integration parameter, either uniformly or according to a weight function.

The method does not require the integrand function to be smooth, we can choose any number of samples we want, and if we decide to request more samples, these can be combined with the existing ones.

Statistically, the error in the Monte Carlo integral estimates will tend to decrease like $N^{-1/2}$. It's good that this number does not depend on the dimension M. But it is nonetheless a remarkably slow rate. Roughly speaking, each new decimal of accuracy requires 100 times as much work.

Since the error goes down relatively slowly, if our initial estimates are poor, it may be a very long time before we achieve a good estimate.



Interpolatory methods sample the integration parameter at selected points, constructing and integrating the interpolating function. If the function is smooth, convergence will be fast, and actually precise if the integrand is a low degree polynomial. Often, a global polynomial is the model; piecewise polynomials are an alternative.

Gauss methods are similar, requiring very specific sampling locations in exchange for squaring the convergence exponent.

For the interpolatory and Gauss rules, the natural approach is to form a **product rule**, using the Cartesian product of the 1D rule taken *M* times. If we approach this problem by asking for a specific precision, then we will specify some number of points *P* for the 1D rule. This implies that our product grid will need $N = P^M$ points.

Even for low precision requests P, we can't go far in dimension M before a product rule becomes unaffordable.



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The guaranteed precision of the product rule comes at a cost which explodes as the dimension M increases.

Each successive entry in a family of quadrature rules integrates precisely more successive monomials of the series *constants, linears, quadratics, cubics, quartics,*

For a given dimension M, if we increase the precision request P, how many monomials must we be able to precisely integrate?

We assume that a quadrature rule can "buy" a monomial at the cost of a function evaluation.



The space of *M*-dimensional *P*-degree polynomials has dimension $\begin{pmatrix} P+M\\ M \end{pmatrix} \approx \frac{M^P}{P!}.$

An *M*-dimensional product of a 1D *P*-point rule requires P^M function evaluations. This means the *M* dimensional product rule can integrate precisely far more monomials than it needs to.

Keep in mind that P is likely to stay relatively small, but M may go to 20, 30, 50 or 100.

As *M* increases, the discrepancy grows. As an extreme example, for M = 100, there are just 101 monomials of degree 1 or less, and 5,151 monomials of degree 2 or less, but a product rule of 2 points would request 2^{100} function evaluations.



The Smolyak approach arises as follows. Let Q be a family of 1D quadrature rules, so that Q(P) can integrate polynomials of degree P exactly. (We'll shortly want to assume that the points used by rules Q(0), Q(1), Q(2) are nested.)

Write any 2D product rule $Q(P_1, P_2) = Q(P_1) \otimes Q(P_2)$

To approximate an integral in 2D up to and including linear terms, we might use the product rule $Q(1,1) = Q(1) \otimes Q(1)$, using 4 points, and precisely integrating any terms involving 1, x, y, or xy.

Smolyak observed that we can get exactly the constants and linears, using just 3 points, using a rule that looks like this:

$$A(L=1,M=2) = Q(1,0) + Q(0,1) - Q(0,0)$$

(The L indicates the level, and M the spatial dimension.)



We're assuming the rules are nested. So we can represent the rules as follows:

Q(0,0) = a * f(0,0) Q(1,0) = b * f(0,0) + c * f(x,0)Q(0,1) = d * f(0,0) + e * f(0,y)

This requires 5 function evaluations before we can combine the three results. Instead, let's combine the coefficients of each function value:

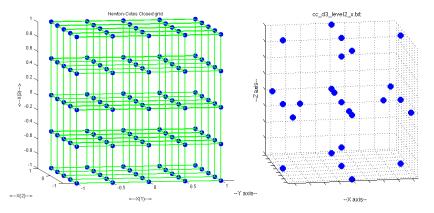
$$\mathcal{A}(L=1,M=2) = (b+d-a) * f(0,0) + c * f(x,0) + e * f(0,y)$$

Now we've guaranteed we only need 3 function evaluations. And we have also written the Smolyak rule as a quadrature rule, that is, coefficients times function values.



SMOLYAK: 5x5x5 Product Versus Smolyak L=2, M=3

125 points in the product rule, 25 points in the Smolyak rule.





 $http://people.sc.fsu.edu/\sim jburkardt/m_src/tensor_grid_display/tensor_grid_display.html$

Smolyak determined a formula combining low-order product rules to achieve a desired precision, avoiding the excessive function evaluation of the standard product rule.

The details of precision, order, efficiency and accuracy vary depending on the underlying 1D quadrature rules.

In most cases, Smolyak matches the product rule while avoiding the crushing explosion in function evaluations.

It's useful now to quote the formal definition of the Smolyak procedure, so you can see that it is framed in an abstract way that allows a lot of flexibility in how it is applied.



We have a family of 1D quadrature rules Q^{ℓ} indexed by ℓ . We form a sparse grid A(L, M) for dimension **M** and level **L**. A(L, M) is a weighted sum of *M*-dimensional product rules

 $\mathcal{Q}^{\ell_1}\otimes \cdots \otimes \mathcal{Q}^{\ell_M}$

The vector $\vec{\ell}$ lists the levels of the component rules used, and $|\vec{\ell}| = \ell_1 + \cdots + \ell_M$ is the sum.

$$\mathcal{A}(L,M) = \sum_{L-M+1 \leq |\vec{\ell}| \leq L} (-1)^{L-|\vec{\ell}|} \begin{pmatrix} M-1 \\ L-|\vec{\ell}| \end{pmatrix} (\mathcal{Q}^{\ell_1} \otimes \cdots \otimes \mathcal{Q}^{\ell_M})$$

Thus, the rule $\mathcal{A}(L, M)$ is a weighted sum of *M*-dimensional product rules whose total level $|\vec{\ell}|$ never exceeds *L*.



We said that the space of **M**-dimensional polynomials of degree **P** or less has dimension $\begin{pmatrix} P+M\\ M \end{pmatrix} \approx \frac{M^P}{P!}$.

For large **M**, a Clenshaw-Curtis Smolyak rule that achieves precision **P** uses $N \approx \frac{(2M)^P}{P!}$ points; we do not see an exponent of *M* in the point count.

For the extreme case of M = 100, a sparse grid based on the Clenshaw Curtis rule can integrate polynomials of total degree 0, 1, 2 or 3 using just 201 points.

Thus, if we are seeking exact integration of polynomials, the Clenshaw-Curtis Smolyak rule uses an optimal number of points (to within a factor 2^{P} that is independent of **M**).



In stochastic problems, the type of random variable used to model uncertainty determines the quadrature rule that should be used to handle that variable. Some common choices include:

Distribution	Domain	Weight	Quadrature	
Uniform	[-1, +1]	1	Gauss-Legendre or	
			Clenshaw-Curtis	
Gaussian	$(-\infty,+\infty)$	$e^{-(x-\alpha)^2/\beta^2}$	Gauss-Hermite	
Gamma	$[0, +\infty)$	$e^{-\alpha x}$	Gauss-Laguerre	
Beta	[-1, +1]	$(1-x)^{lpha}(1+x)^{eta}$	Gauss-Jacobi	

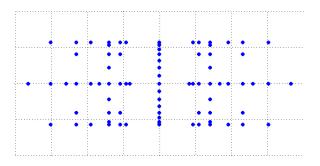
Any of these rules can be used in the sparse grid formulation.



SMOLYAK: Mixed Products

The Smolyak procedure allows each dimension to be treated independently in whatever manner is appropriate for that variable.

The sparse grid technique simply combines the given rules, without requiring that they involve the same 1D domain or weight function.





Level 4 Rule, Hermite in X, Clenshaw Curtis in Y http://people.sc.fsu.edu/~jburkardt/m_src/sparse_grid_mixed_dataset/sparse_grid_mixed_dataset.html Refer to **Novak and Ritter*** for details about the construction of rules based on the Clenshaw Curtis rule, error bounds for the approximation of integrals of smooth functions, estimates for the growth in the number of function evaluations with increasing level or dimension, and a proof that the rules are exact for all monomials up to a given degree.

The authors also indicate conditions under which other 1D quadrature rules can be used to construct Smolyak rules of comparable precision.



^{*}High dimensional integration of smooth functions over cubes, Numerische Mathematik, volume 75, pages 79-97, 1996.

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The Smolyak definition does not specify the properties of the 1D quadrature rules used to form product rules.

A popular choice uses the 1D Clenshaw Curtis rule (based at the zeros of the Chebyshev polynomials). A nested sequence of rules is selected, of orders 1, 3, 5, 9, 17, 33, ... so the rule of one level includes the points of the previous one. This is the approach in Novak and Ritter.

Nesting keeps our point count low by reusing points; on the other hand, rules of successive levels double in size, which is a potential (but fixable) pitfall.



ASPECTS: Computing X and W for Clenshaw Curtis

MATLAB code to compute a Clenshaw Curtis rule of order N:

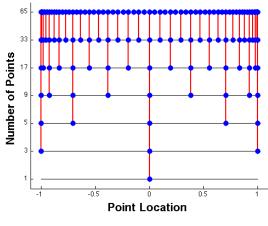
```
theta(1:n) = (n - 1 : -1 : 0) * pi / (n - 1):
x(1:n) = cos ( theta(1:n) );
w(1:n) = 1.0:
for i = 1 : n
 for i = 1; floor ((n - 1)/2)
   if (2 * i = (n - 1))
     b = 1.0;
   else
     b = 2.0;
    end
   w(i) = w(i) - b * cos ( 2 * j * theta(i) ) / ( 4 * j * j - 1 );
 end
end
w(1)
                w(1) / (n - 1):
        =
w(2:n-1) = 2.0 * w(2:n-1) / (n - 1);
w(n)
        =
                w(n)
                        / (n - 1);
```

http://people.sc.fsu.edu/~jburkardt/m_src/quadrule_fast/quadrule_fast.m



ASPECTS: We Choose a Nested Clenshaw Curtis Sequence

Each grid includes all points from the previous one. (However, the size of the grid grows more rapidly than required.)



Nested Clenshaw Curtis Sequence



http://people.sc.fsu.edu/~jburkardt/m_src/nested_sequence_display/nested_sequence_display.html

We can index our rules by level L, writing CC(0) for the the first rule, counting the points in the rule as the order O, and denoting the polynomial precision by P.

Rule	CC(0)	CC(1)	CC(2)	CC(3)	CC(4)	CC(5)	
L	0	1	2	3	4	5	
0	1	3	5	9	17	33	$2^{L-1} + 1$
Р	1	3	5	9	17	33	P=O

Because our 1D rules have odd order, instead of P=O-1 we will actually have P=O, that is, a 1 point rule gets constants **and** linears, a three point rule picks up cubics, and so on.

 $http://people.sc.fsu.edu/{\sim}jburkardt/cpp_src/sparse_grid_cc/sparse_grid_cc.html$

ASPECTS: The First Three CC Rules

The CC-based Smolyak rules $\mathcal{A}(L = 0/1/2, M = 2)$ are:

 $\mathcal{A}(0,2) = CC(0) \otimes CC(0)$

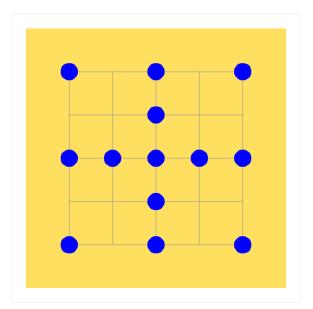
 $egin{aligned} \mathcal{A}(1,2) = & CC(1) \otimes CC(0) \ & + & CC(0) \otimes CC(1) \ & - & CC(0) \otimes CC(0) \end{aligned}$

$$\begin{aligned} \mathcal{A}(2,2) = & CC(2) \otimes CC(0) \\ & + & CC(1) \otimes CC(1) \\ & + & CC(0) \otimes CC(2) \\ & - & CC(1) \otimes CC(0) \\ & - & CC(0) \otimes CC(1) \end{aligned}$$

(For higher dimensions, we don't just have +1 and -1 as coefficients, but they will still be combinatorial coefficients.)



ASPECTS: The Quadrature Points of $\mathcal{A}(2,2)$

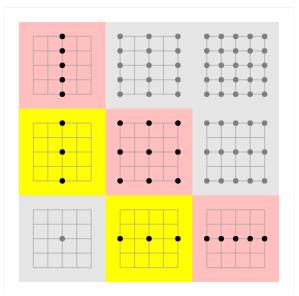




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ASPECTS: A(2,2) = 2x0 + 1x1 + 0x2 - 1x0 - 0x1

It helps to see the underlying product rules that were combined:





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The picture correctly suggests that the Smolyak combination of the 5 lower order grids is (essentially) as precise as the 5x5 product grid in the upper right corner.

But the Smolyak grid $\mathcal{A}(2,2)$ uses 13 points, the product grid $C(5)\otimes C(5)$ uses 25.

The Smolyak definition chooses a collection of lower order product grids that capture the information necessary to approximate all the monomials of interest, and then combines this information correctly to produce a good integral estimate.

(Because of the nesting choice we made, the precision results are not so neat for higher levels...)



Novak and Ritter show that for sparse grids based on the 1D Clenshaw Curtis rule, the precision is related to the level by:

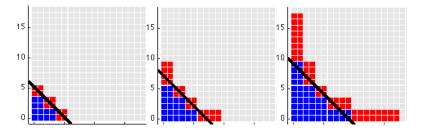
$$P = 2 * L + 1$$

The first sparse grid (L=0) picks up constants and linears, the second adds quadratics and cubics, and so on. This means that our family of sparse grids has a predictable precision behavior based on the level, regardless of dimension.

If our CC-based sparse grids were perfectly efficient, the level L rule would pick up exactly the monomials up to precision P and no more. A precision graph will show that we do better than a product rule, but still have some inefficiency.



Here are the precisions of $\mathcal{A}(2,2), \mathcal{A}(3,2), \mathcal{A}(4,2)$. The diagonal black line "fences off" the monomials of degree 2*L+1, which the sparse grid must integrate precisely. With increasing level, the rules spill over the fence, suggesting some inefficiency.



(The red squares are monomials just added on this level.)



ASPECTS: Reducing the Nesting Overhead

Our sparse grid of level L will have the precision P = 2 * L + 1 as long as our family of 1D rules has at least that same precision sequence:

L		0	1	2	3	4	5	6	7	8	9	10
Ρ	minimal:	1	3	5	7	9	11	13	15	17	19	21

But our 1D rules double the number of points with each level, in order to effect nesting...which we want to control the point count.

P supplied: 1 3 5 9 17 33 65 129 255 513 1025...

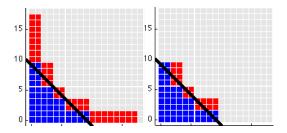
We could keep nesting but reduce the cost of doubling, by reusing a 1D rule if it satisfies the precision requirement.

P suggested: 1 3 5 9 9 17 17 17 17 33 33...

http://people.sc.fsu.edu/~jburkardt/cpp_src/sgmg/sgmg.html



Using the slower growth strategy preserves our overall precision requirement while cutting down on the order of the finest grid, and hence the number of function evaluations.



Both rules achieve precision P = 2 * L + 1 = 9; the standard rule uses 65 points, the slow rule 49.



The standard formulation of the Smolyak formula treats each dimension equally.

In the precision diagrams, this corresponds to using a diagonal line that indicates the total polynomial precision we expect on each successive rule.

We might now in advance, or discover adaptively, that some coordinates are inactive, and that some are relatively more active than others. It is easy to modify the Smolyak formula so that, with each successive rule adds significantly more precision (more monomials) in the most active dimensions.

In this case, our precision diagram (in 2D) would involve a sequence of nondiagonal lines whose slope corresponds to the level of anisotropy.

 $http://people.sc.fsu.edu/{\sim}jburkardt/cpp_src/sgmga/sgmga.html$



It turns out that the Smolyak formula can also handle very general cases suitable for adaptivity.

A given 2D precision diagram is a monotonically decreasing "cityscape" of squares representing monomials we have captured.

Each "corner" of the cityscape represents a monomial we didn't capture, but which could be captured by including a single new product rule to the Smolyak formula.

Thus, an adaptive approach could start with a given precision diagram, then consider each corner and compute its potential new contribution, and take those which rise above a given tolerance in magnitude.

Naturally, the search approach itself becomes expensive as the dimension gets very high.

 $http://people.sc.fsu.edu/{\sim}jburkardt/cpp_src/sandia_sgmgg/sandia_sgmgg.html$



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We estimate the integral of f(x) over an *M*-dimensional hypercube $[a, b]^M$ using *N* points or function evaluations.

The "cost" of the estimate is N. If an estimate is not satisfactory, it's important to know how fast the error is likely to go down if we increase N and try again.

For the Monte Carlo method, the rate of error decay is known to be $O(\frac{1}{\sqrt{N}})$. The rate is independent of spatial dimension M, and essentially independent of the smoothness of f(x).

Compare Monte Carlo and sparse grid values of N and accuracy.



NUMERICAL: Sparse Grids May Require Smoothness!

Let f(x) be the characteristic function of the unit ball in 6D:

Ν	SG Estimate	SG Error	:	MC Estimate	MC Error
1	4.000	1.167	:	0.00000	5.16771
13	64.000	58.832	:	0.00000	5.16771
85	-42.667	-47.834	:	3.01176	2.15595
389	-118.519	-123.686	:	4.77121	0.39650
1457	148.250	143.082	:	5.16771	0.01555
4865	-24.682	-29.850	:	5.41994	0.25226

Can you see why negative estimates are possible for the sparse grid, even though the integrand is never negative?

Sparse grids based on global polynomials require smooth integrands; in this case, a sparse grid using piecewise polynomials would have succeeded.



 $\label{eq:http://people.sc.fsu.edu/~jburkardt/m_src/quadrature_test/quadrature_test.html, problem \#18 \\ \mbox{http://people.sc.fsu.edu/~jburkardt/m_src/ball_volume_monte_carlo/ball_volume_monte_carlo.html} \label{eq:http://people.sc.fsu.edu/}$

Monte Carlo doesn't diverge, but look how hard we have to work to get three places of accuracy for the characteristic function of the unit ball in 6D.

N	MC Estimate	MC Error	
1	0.00000	5.16771	
32	6.00000	0.83228	
1,024	4.81250	0.35521	
32,768	5.39063	0.22291	
1,048,576	5.18042	0.01271	
33,554,432	5.16849	0.00077	
∞	5.16771	0.00000	

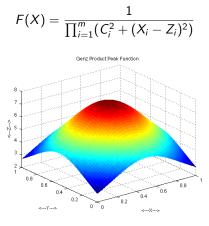
Should we want one more digit of accuracy, we can expect to need 100 times as many points \approx 3.3 billion points.

http://people.sc.fsu.edu/~jburkardt/m_src/ball_volume_monte_carlo/ball_volume_monte_carlo.html



NUMERICAL: Genz Product Peak Test in 6D

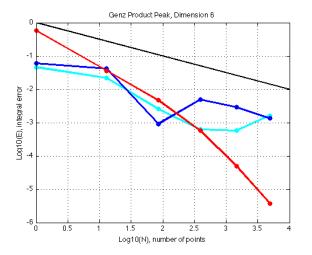
Alan Genz provided six high dimensional test integrals; The **product peak function** is defined on the unit hypercube, with given C and Z vectors, and is smooth:





NUMERICAL: Genz Product Peak Test in 6D

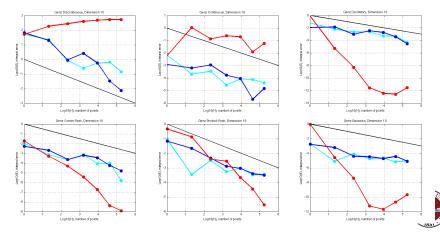
Red: Sparse grid estimate Blue & Cyan: MC estimates Black: Expected MC Rate of Decrease





NUMERICAL: Genz Tests in 10D

Discontinuous, Continuous, Oscillatory Corner Peak, Product Peak, Gaussian (sparse grid estimate in red)



Let's return to the stochastic Poisson equation considered earlier:

$$-\nabla \cdot (a(x,y;\omega) \nabla u(x,y;\omega)) = f(x,y)$$

Our integration problem seeks the expected value of $u(x, y; \omega)$, assuming we have a probabilistic model for the stochastic influence.

Monte Carlo: select a random set of parameters ω according to $pr(\omega)$, solve the Poisson equation for u, and average.

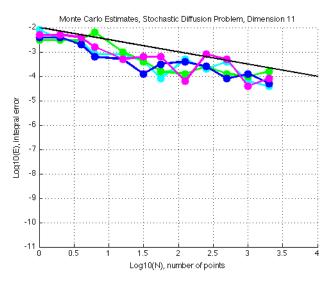
Sparse grid: choose a level, defining a grid of ω values, solve the Poisson equation for u, multiply by the probability, and take a weighted average.

Clayton Webster, Sparse grid stochastic collocation techniques for the numerical solution of partial differential equations with random input data, PhD Thesis, Florida State University, 2007.



NUMERICAL: Four Monte Carlo Estimates

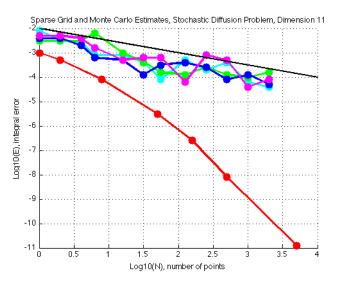
The black line is the Monte Carlo trend.





NUMERICAL: Sparse Grid Versus Monte Carlo

The sparse grid estimates converge rapidly.





For the stochastic diffusion problem, $u(x, y; \omega)$ has a very smooth dependence on the perturbations in ω .

For this reason, a sparse grid can sample the solution for a small set of perturbations ω and accurately estimate the expected value.

If we had a code to solve the original Poisson equation for a given conductivity field, the sparse grid procedure simply needs to call that unmodified code with different conductivities.

This is why sparse grids are called a nonintrusive method. Other procedures for dealing with uncertain or stochastic influences may require extensive changes, new variables, and a larger coupled system to solve.



- Introduction
- Integrals Arising from Stochastic Problems
- Smolyak's Sparse Grid Definition
- Aspects of Sparse Grids
- Numerical Examples
- Conclusion



In general, estimating the integral of an arbitrary function in a high dimensional space is an intractable problem.

Nonetheless, we are able to estimate the high dimensional integrals associated with stochastic problems and uncertainty quantification when the mathematical modeling of uncertainty results in integrands that are <u>smooth</u>, or <u>localized</u>, or can be well approximated by certain easily integrated density functions.

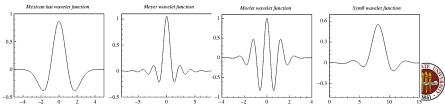
Sparse grids are a powerful tool for treating such special integrands, whose estimated integrals can be pulled out of an otherwise impenetrable high dimensional jungle!



The sparse grid construction can be modified so that more precision is gained in certain coordinate directions. Essentially, the diagonal lines in the precision diagram become slanted. This coordinate preference can be done beforehand, or can be modified during execution.

We are free to plug into Smolyak's formula any sequence of 1D quadrature rules, including Gauss-Legendre, composite rules, or a hierarchical family based on piecewise linear hat functions.

Clayton Webster is looking at a method that employs the unusual choice of **wavelets** to carry out the 1D approximations.



The sparse grid procedure I described assumes smooth functions. For a function with a "mild" discontinuity , an adaptive procedure can be devised to detect the problem, and patch up the 1D quadrature schemes. The hierarchical linear basis is a natural way to do this.

Let f(x) be defined on $[0,1]^N \to [0,1]$ as the product of N 1D step functions, each jumping at a random location ξ_i . Because the function is separable, I can locate the problem areas by looking at each 1D projection, and refining the 1D quadrature rule in dimension *i* near ξ_i .

But now let g(x) be f(Q * x), where Q is an orthogonal matrix which rotates the coordinates. The discontinuity lines do not align with the coordinate axes, and the location problem becomes much harder.

And this is nowhere near as bad as the problem could become!



SPINTERP, a MATLAB program, by Andreas Klimke, is a great way to explore the power of sparse grids.

```
x = spgrid (l, m)
```

returns the points of a sparse grid of level $\boldsymbol{\mathsf{L}}$ in dimension $\boldsymbol{\mathsf{M}}.$

To estimate an integral:

```
z = spvals ( @fun, m )
q = spquad ( z )
```

SPINTERP can also interpolate and optimize using sparse grids.

http://www.ians.uni-stuttgart.de/spinterp/



SPARSE_GRID_HW, a MATLAB library, by Florian Heiss and Viktor Winschel, is both a calculator to use, and a toolbox you can learn from and build with.

- It includes 1D quadrature families of Legendre and Hermite types.
- It includes Kronrod-Patterson variants of these families for slower growth.
- It makes it easy for the user to define a new family, simply by supplying a table of sizes, points and weights for the 1D rules.
- I was able to make C, C++ and FORTRAN90 versions as well.

 $http://people.sc.fsu.edu/{\sim}jburkardt/m_src/sparse_grid_hw/sparse_grid_hw.html$



CONCLUSION: References

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